

Quark and gluon contributions to the spin and momentum of the nucleon from lattice QCD

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POETIC7, 15 November 2016,
Temple University,
Philadelphia



THE CYPRUS
INSTITUTE

CaSToRC

★ Results here by Cyprus – DESY – Temple

C. Alexandrou, S. Bacchio, M. Constantinou, K.
Hadjiyiannakou, K. Jansen, C. Kallidonis, F. Steffens, A.
Vaquero, C. Wiese

★ ETMC collaboration

Cyprus (U. Cyprus and Cyprus Inst.)
Germany (Berlin/Zeuthen, Bonn, Frankfurt)
Italy (Rome I,II,III, Trento)
Poland (Poznan)
Switzerland (Bern)
UK (Liverpool)

Outline

★ Short introduction to lattice calculations

- Challenges and current landscape

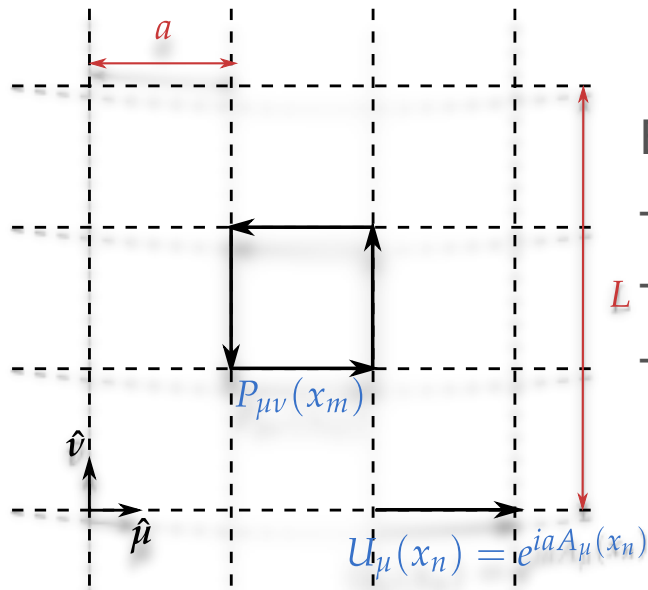
★ Nucleon spin on the lattice

- Axial charges
- Moments of PDFs
- Spin decomposition
- Gluon moment

★ Summary and outlook

Lattice QCD – *ab initio* simulation of QCD

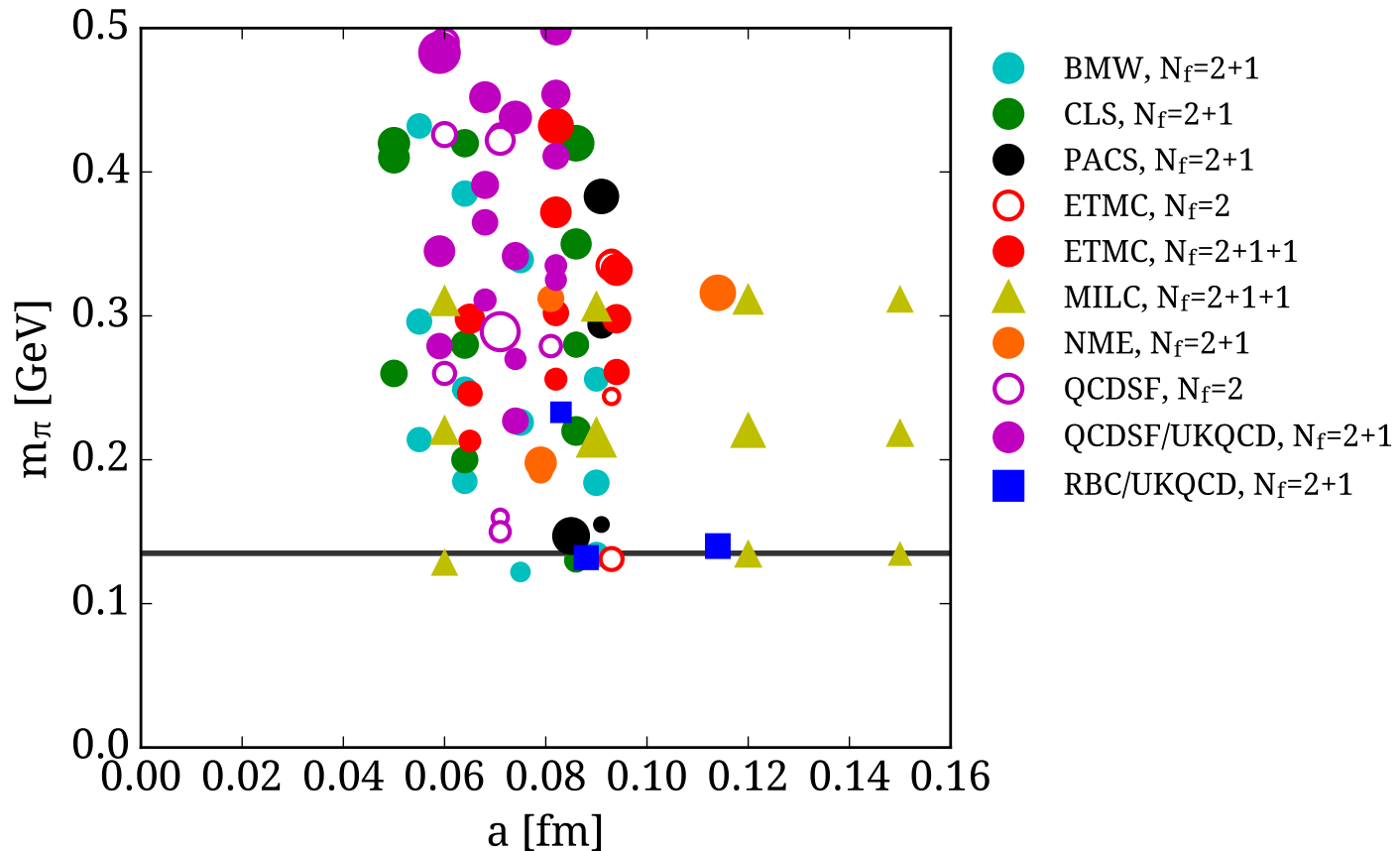
- Freedom in choice of:
 - quark masses (*heavier is cheaper*)
 - lattice spacing a (*larger is cheaper*)
 - lattice volume $L^3 \times T$ (*smaller is cheaper*)
- Choice of discretisation scheme
e.g. **Clover, Twisted Mass, Staggered, Overlap, Domain Wall**
Trade – offs and advantages for each differ



Eventually, **all** schemes must agree:

- At the continuum limit: $a \rightarrow 0$
- At infinite volume limit $L \rightarrow \infty$
- At physical quark mass

Simulations landscape

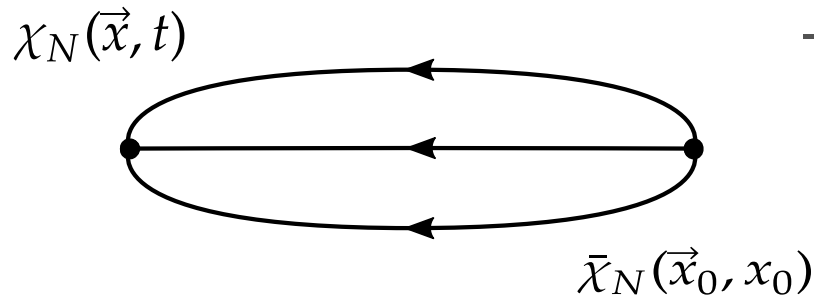
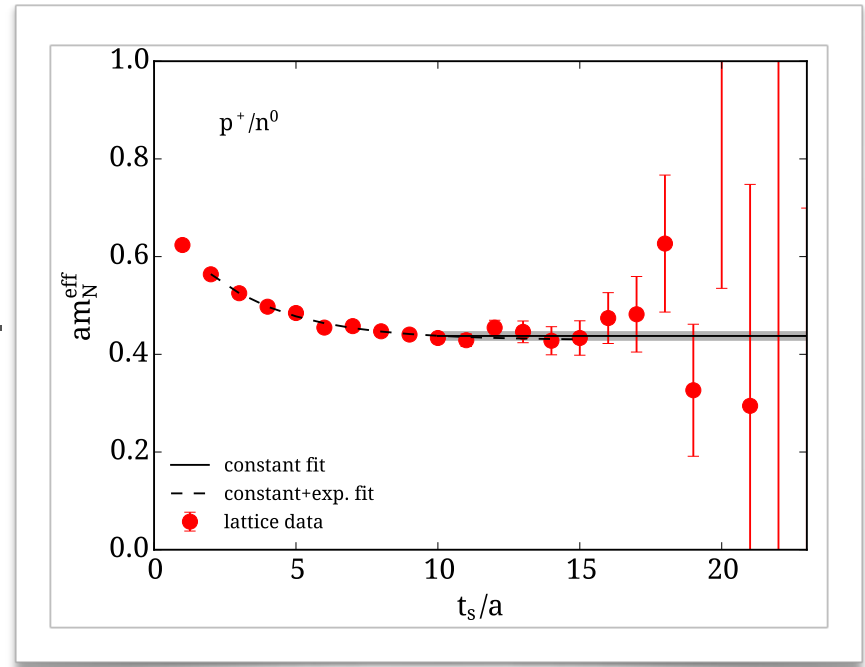


Selected lattice simulation points used for hadron structure

- Multiple collabs. simulating at physical pion mass
- Size of points indicates $m_\pi L$

Sources of uncertainty

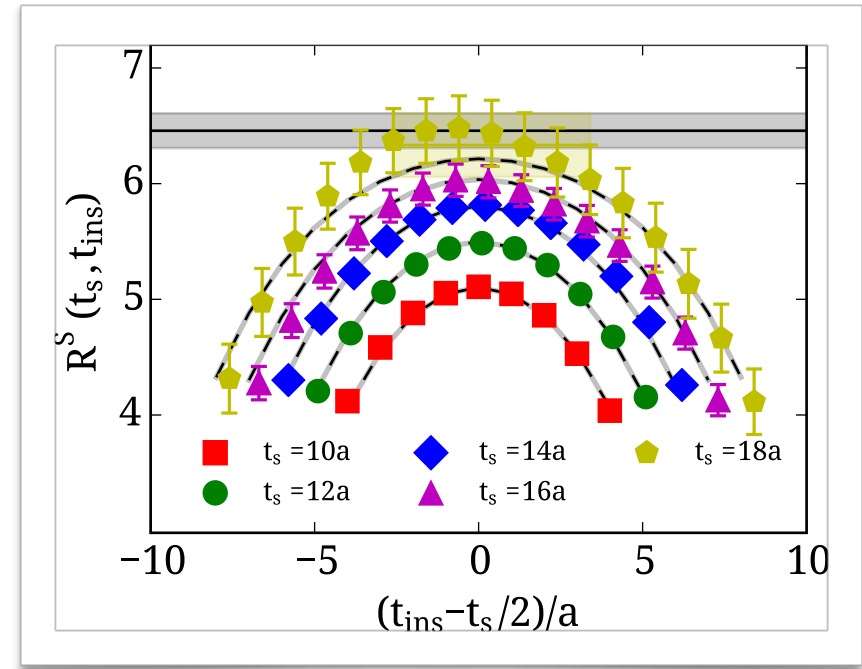
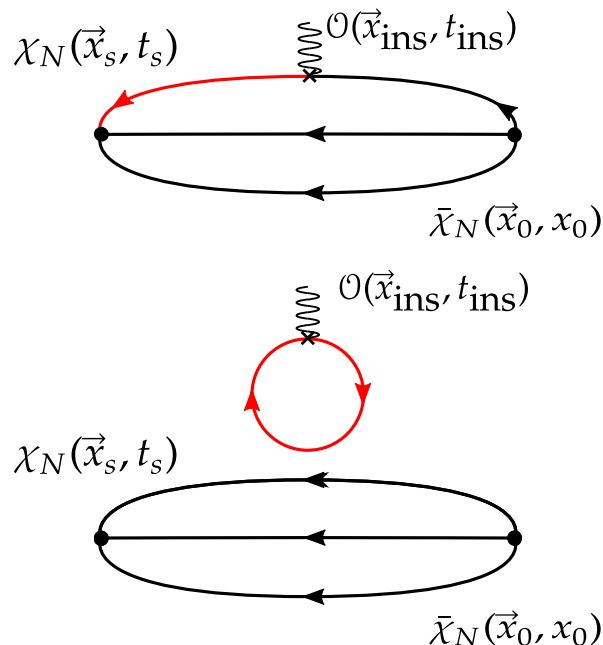
- Statistical error: $\frac{1}{\sqrt{N}}$, with MC samples
- Correlation functions: exponentially decay with time-separation
- Disconnected contributions: stochastic error



- Systematic uncertainties
 - Extrapolations
 a, L, m_π
 - Contamination from higher energy states

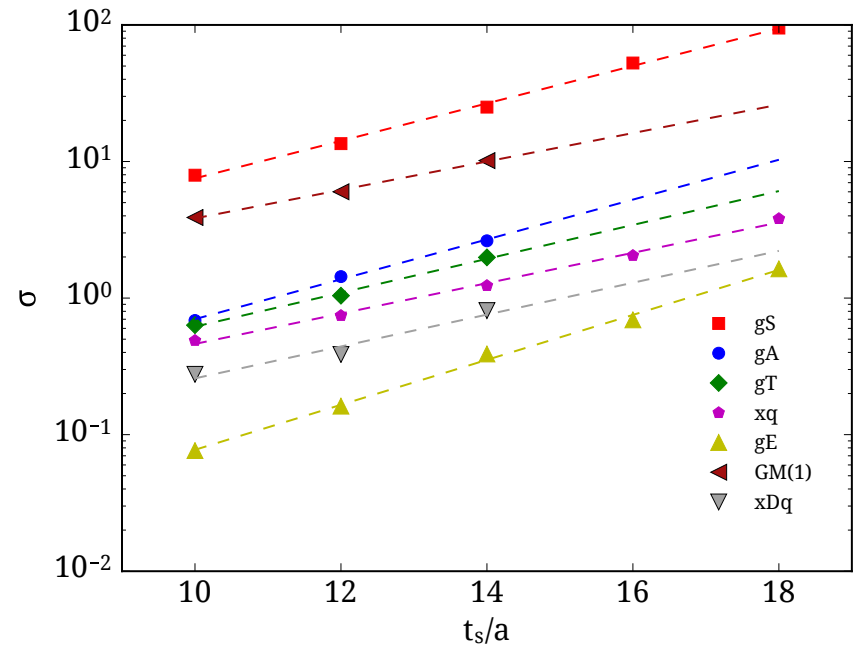
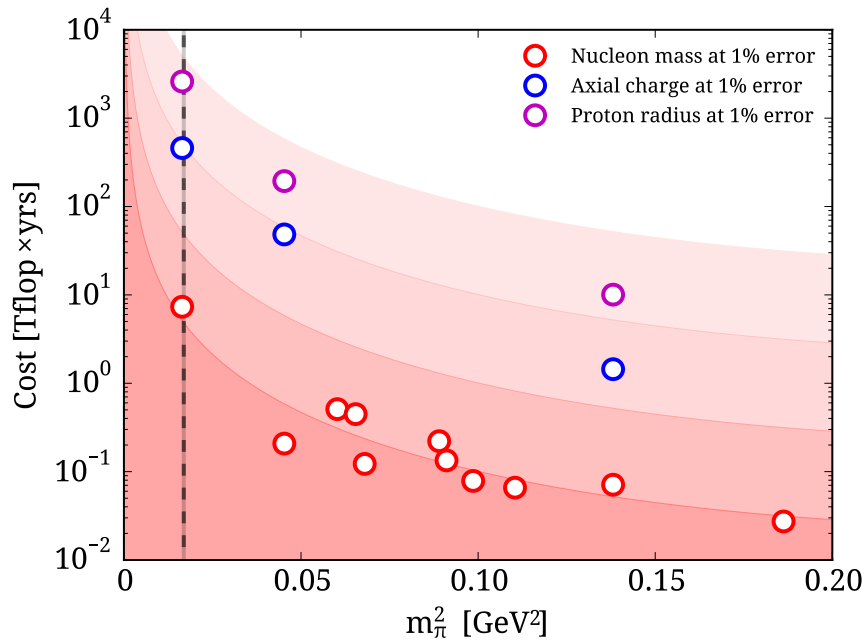
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- Systematic uncertainties
 - Extrapolations a, L, m_π
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Multi-petascale to exa-scale requirements

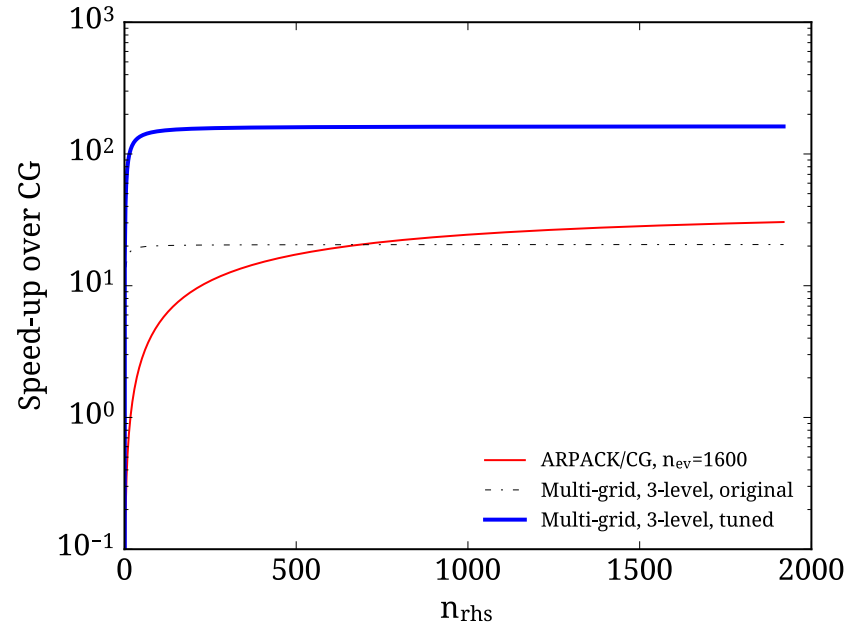
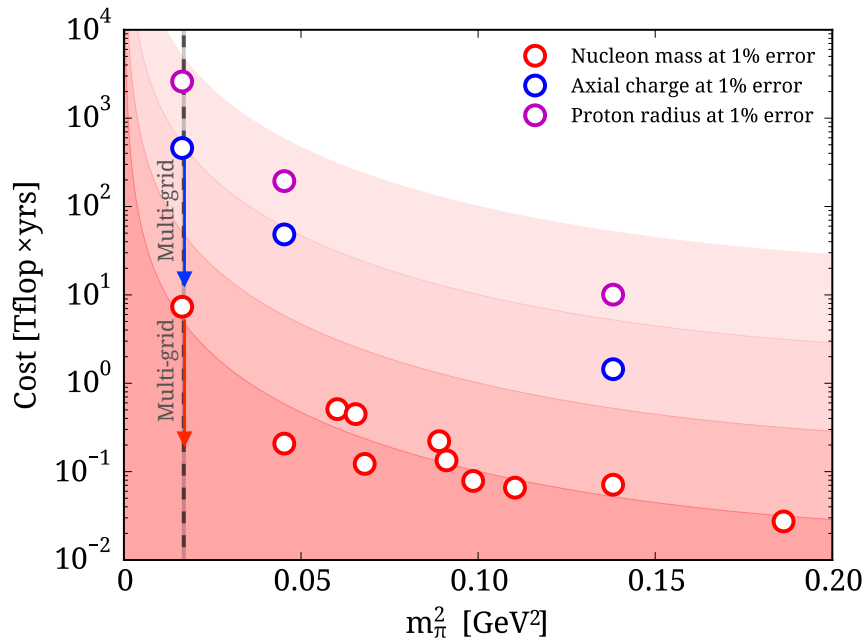


Indicative computer time requirements for nucleon structure



Increased time separations required for suppression of excited states at physical point

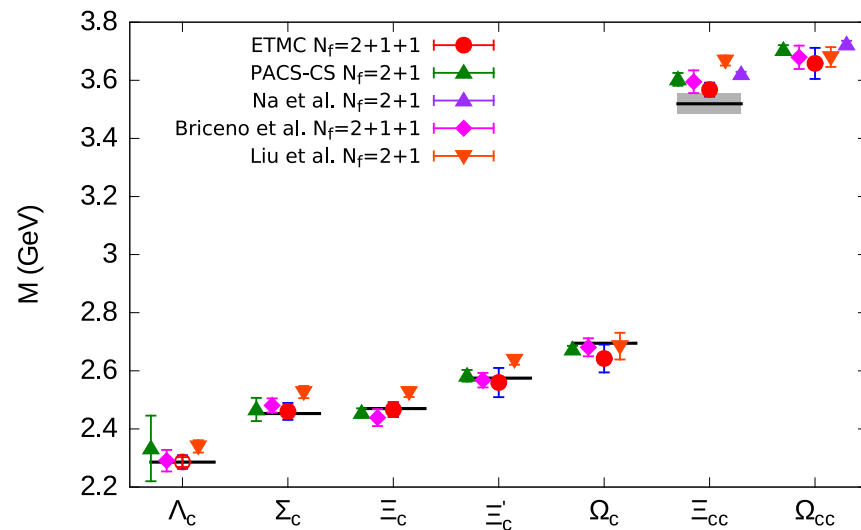
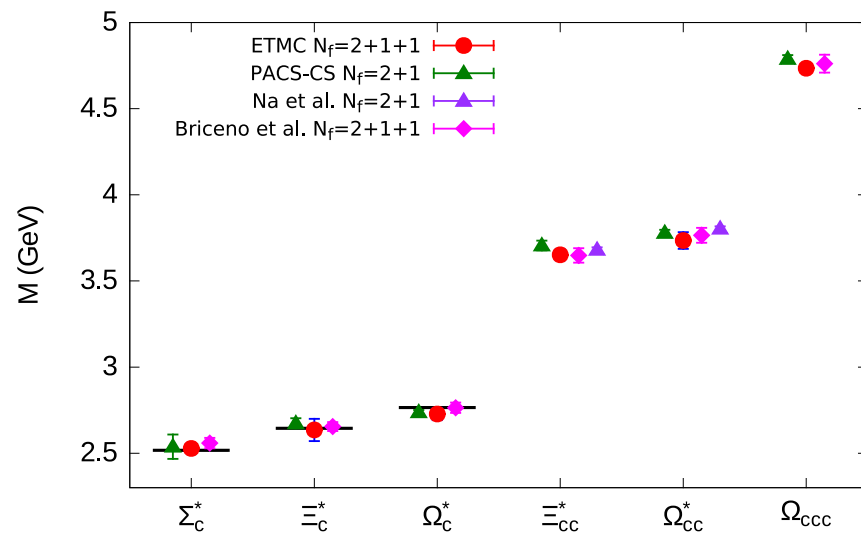
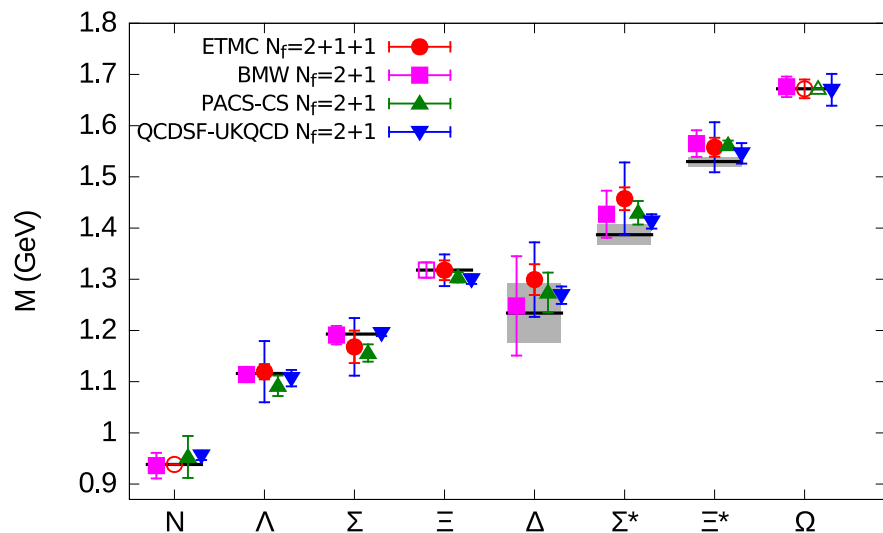
Multi-petascale to exa-scale requirements



Indicative computer time requirements for nucleon structure



Multi-grid/deflation methods for large speed-ups, especially at physical point



Reproduction of light baryon masses

- Agreement between lattice discretisation schemes
- Reproduction of experiment

Prediction of yet-to-be-observed charmed baryons

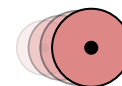
- Confidence through agreement between lattice schemes
- Nucleon structure...

Nucleon structure on the lattice

- Lattice: moments of GPDs are readily accessible

Unpolarised

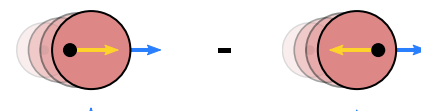
$$\mathcal{O}_{\boxed{V}}^{\mu\mu_1\mu_2\ldots\mu_n} = \bar{\psi} \boxed{\gamma}^{\mu} iD^{\mu_1} iD^{\mu_2} \ldots iD^{\mu_n} \psi$$



$$\langle 1 \rangle_{u-d} = g_V, \quad \langle x \rangle_{u-d}, \quad \dots$$

Helicity

$$\mathcal{O}_{\boxed{A}}^{\mu\mu_1\mu_2\ldots\mu_n} = \bar{\psi} \boxed{\gamma_5 \gamma}^{\mu} iD^{\mu_1} iD^{\mu_2} \ldots iD^{\mu_n} \psi$$



$$\langle 1 \rangle_{\Delta u - \Delta d} = g_A, \quad \langle x \rangle_{\Delta u - \Delta d}, \quad \dots$$

Transverse

$$\mathcal{O}_{\boxed{T}}^{\nu\mu\mu_1\mu_2\ldots\mu_n} = \bar{\psi} \boxed{\sigma}^{\nu} iD^{\mu_1} iD^{\mu_2} \ldots iD^{\mu_n} \psi$$

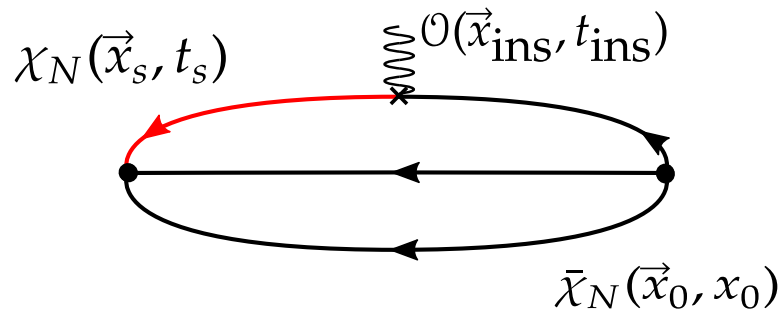


$$\langle 1 \rangle_{\delta u - \delta d} = g_T, \quad \langle x \rangle_{\delta u - \delta d}, \quad \dots$$

Axial charge

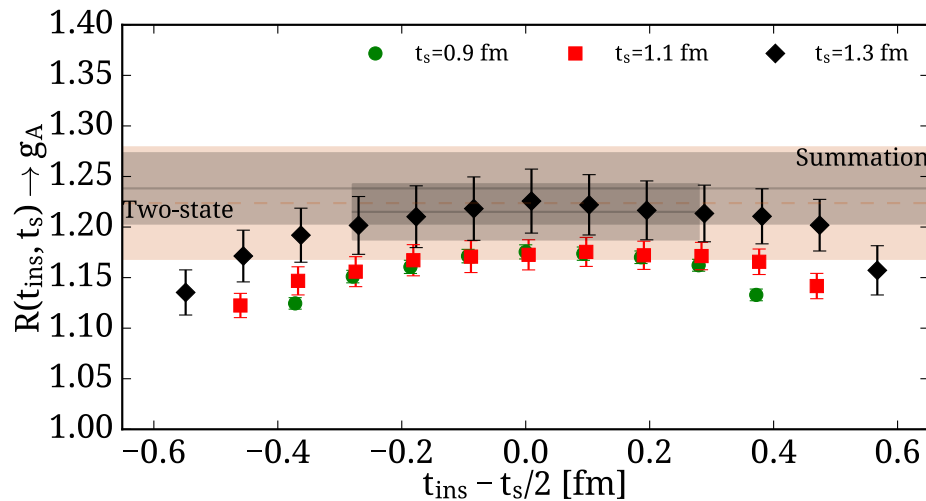
Isovector matrix element:

$$\langle N | \bar{u} \gamma_5 \gamma_k u - \bar{d} \gamma_5 \gamma_k d | N \rangle \rightarrow g_A$$



Need simultaneously:

$$t_{ins} \gg, \quad t_s - t_{ins} \gg$$



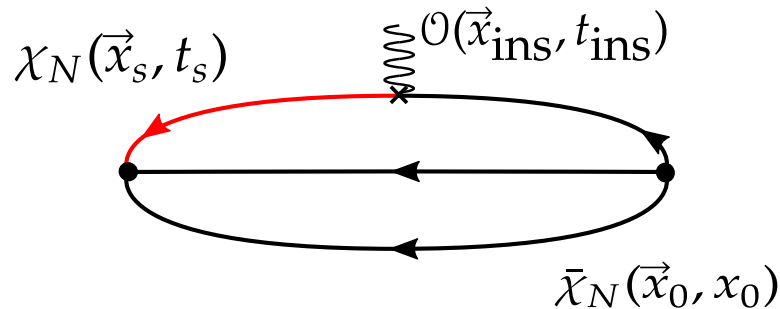
- Direct fit to single t_s
- Fit to multiple t_s including an excited state
- Fit to sum over t_{ins} as function of t_s

$m_\pi = 135$ MeV, Twisted Mass: Phys.Rev. D92 (2015) 114513

Axial charge

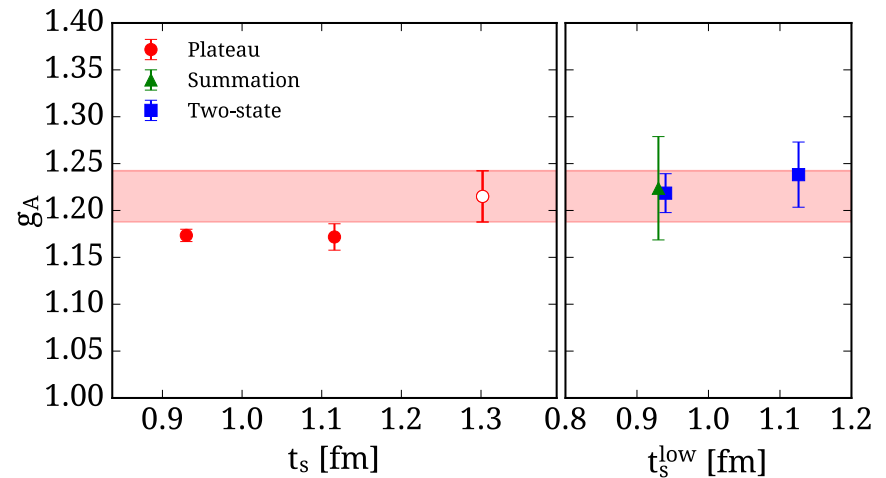
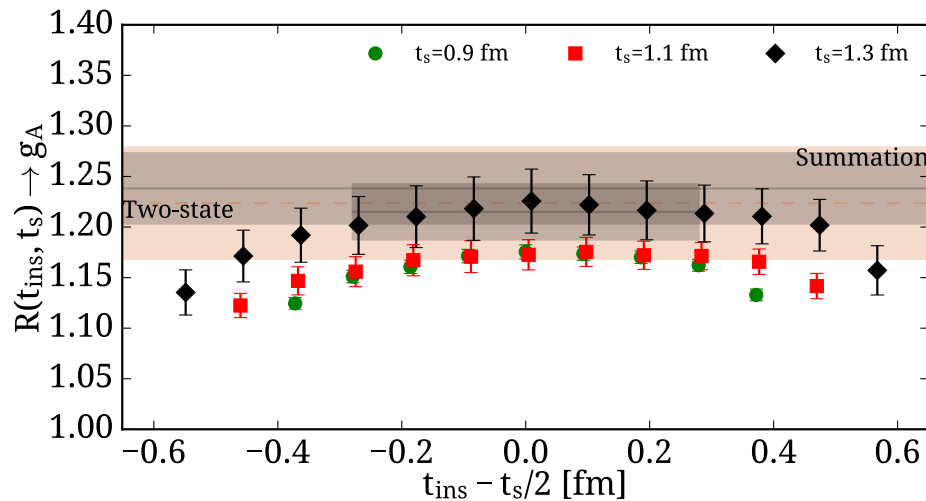
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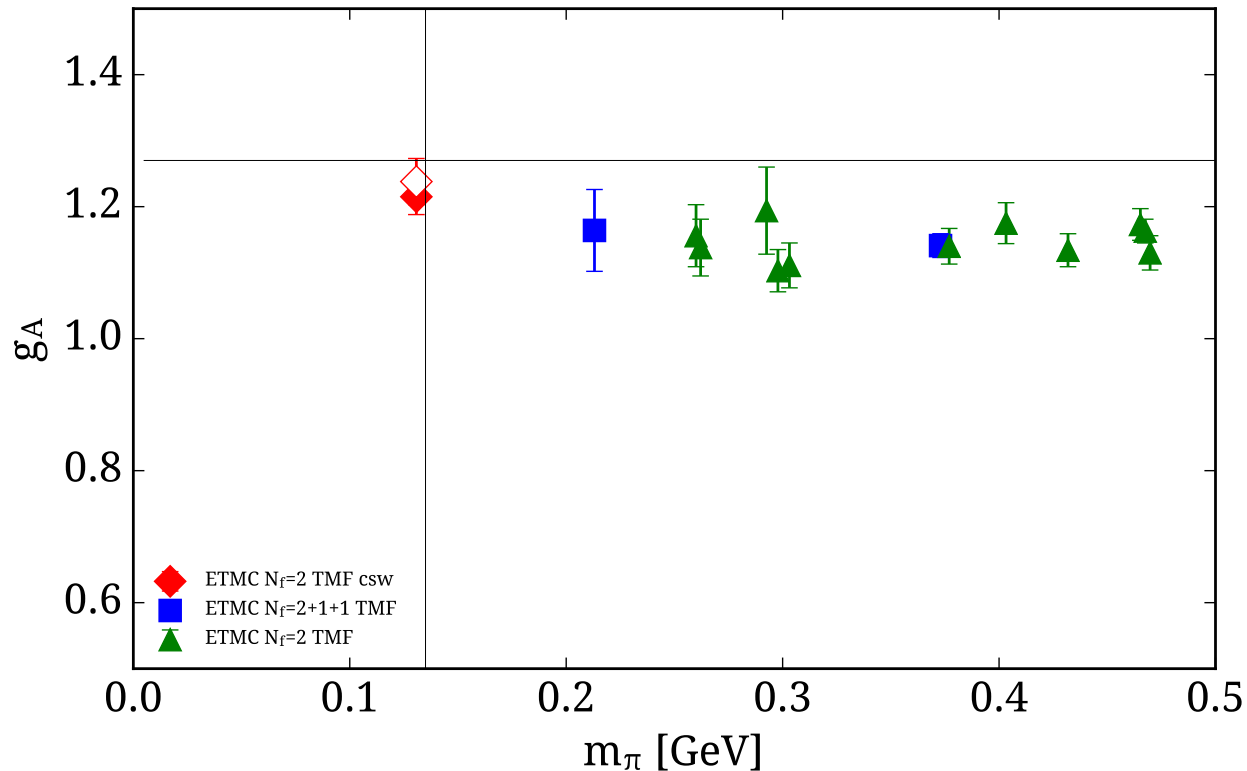
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$m_\pi = 135$ MeV, Twisted Mass: Phys.Rev. D92 (2015) 114513

Axial charge

$$\langle N | \bar{u} \gamma_5 \gamma_k u - \bar{d} \gamma_5 \gamma_k d | N \rangle \rightarrow g_A$$

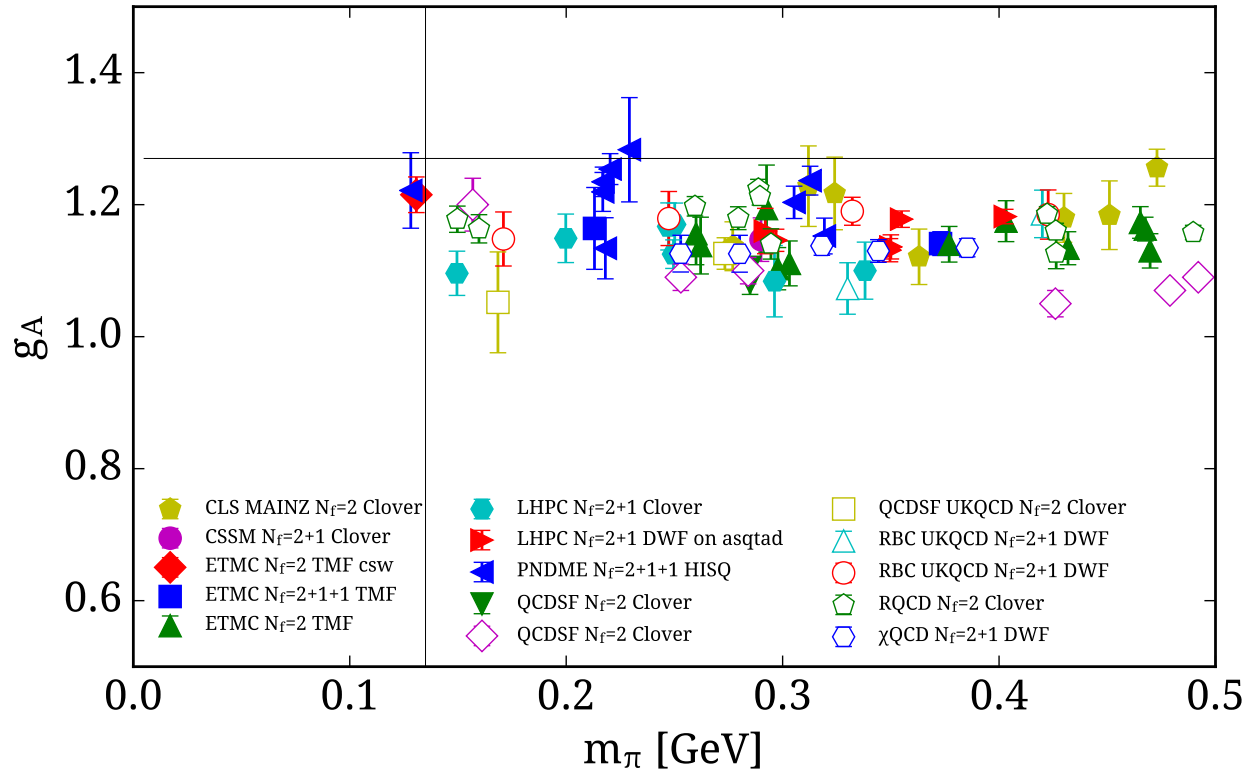


Axial charge

- Agreement towards experiment
- Simulations very close to or at the physical quark mass
- At physical point: statistics of $\sim 10^4$, up to 1.3 fm separation

Axial charge

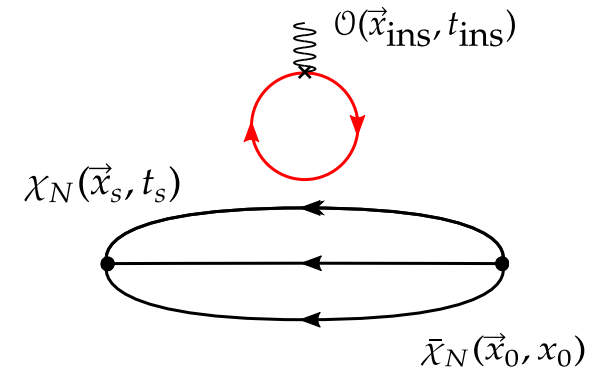
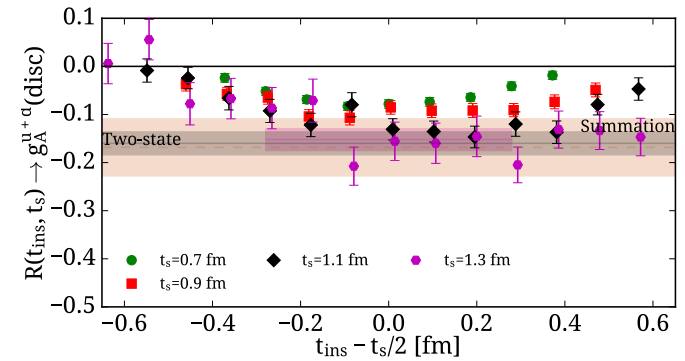
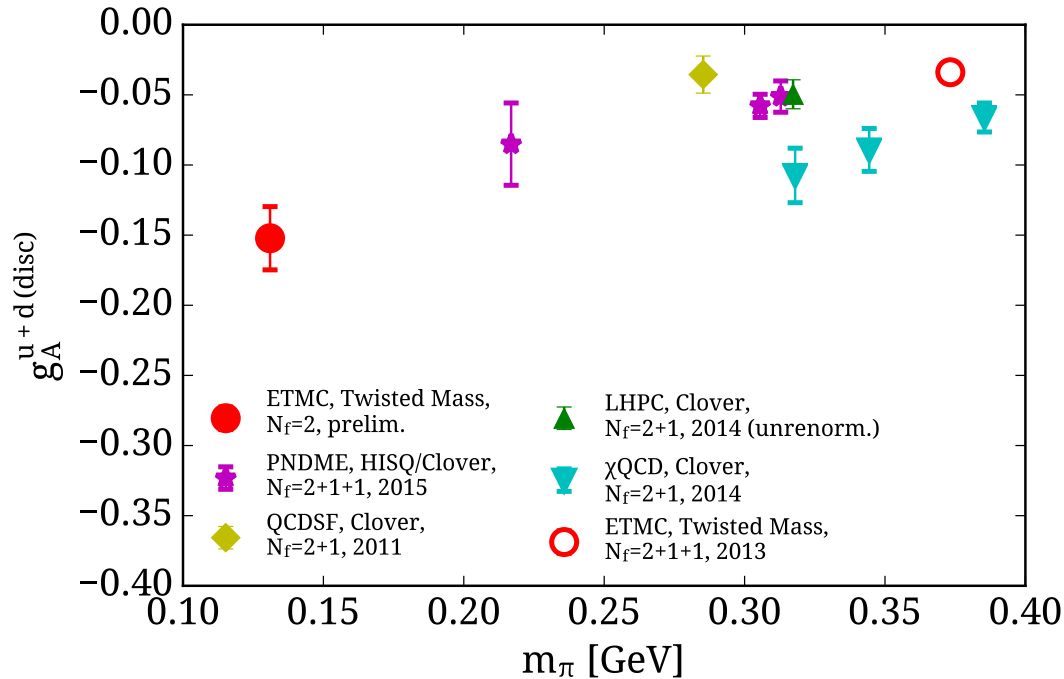
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Axial charge

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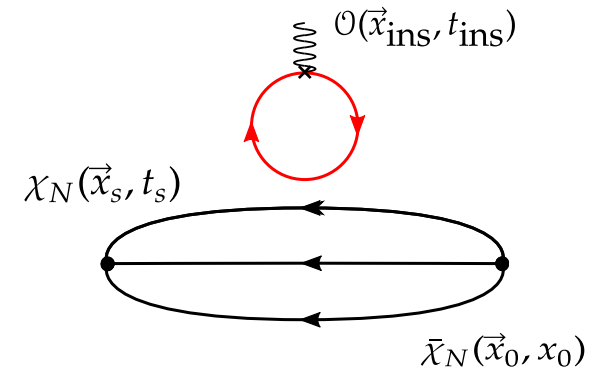
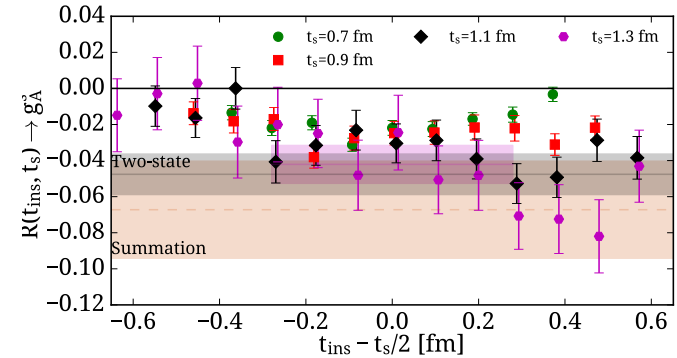
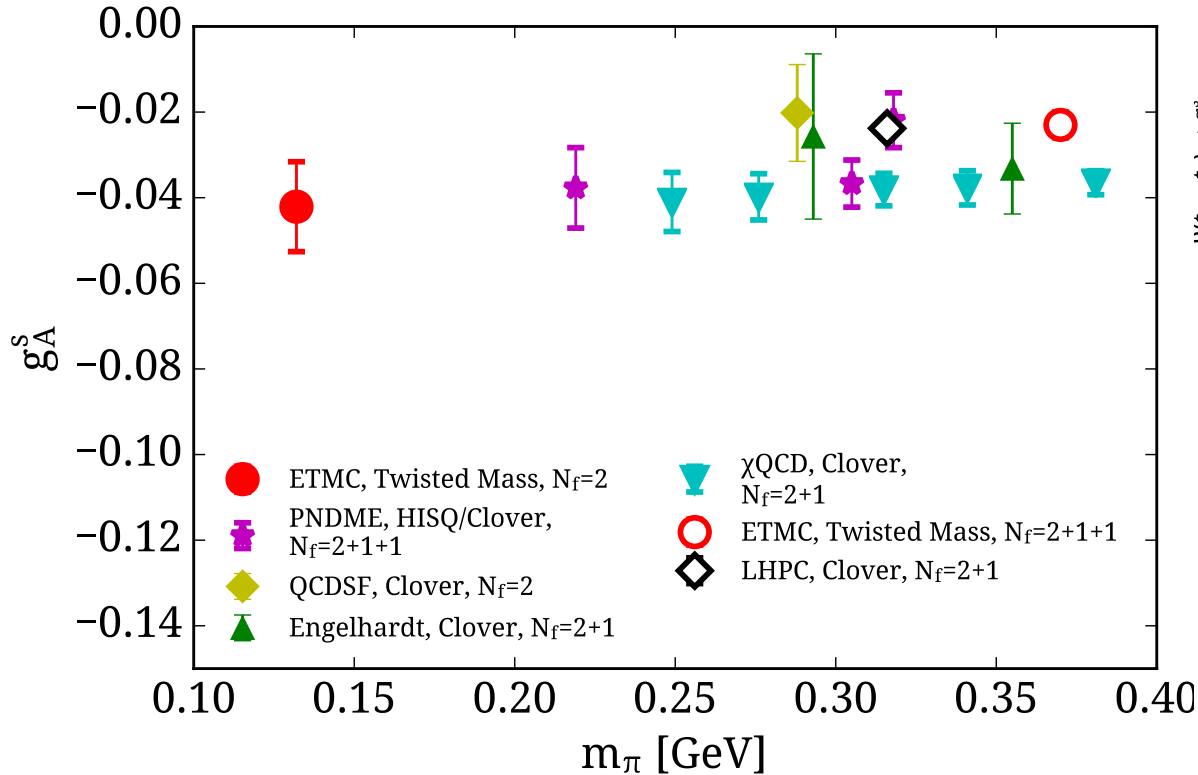
$$\bar{u}\gamma_5\gamma_\mu u + \bar{d}\gamma_5\gamma_\mu d$$



Axial charge — light disconnected

- Required for individual u- and d- contributions
- Requires dedicated calculations for “disconnected quark loop”
- Large statistical fluctuations in correlation functions
- $>10^5$ statistics at all available separations
- Sign is negative: brings connected result down
- About 10% of connected value

$$\bar{s}\gamma_5\gamma_\mu s$$



Axial charge – strange contribution

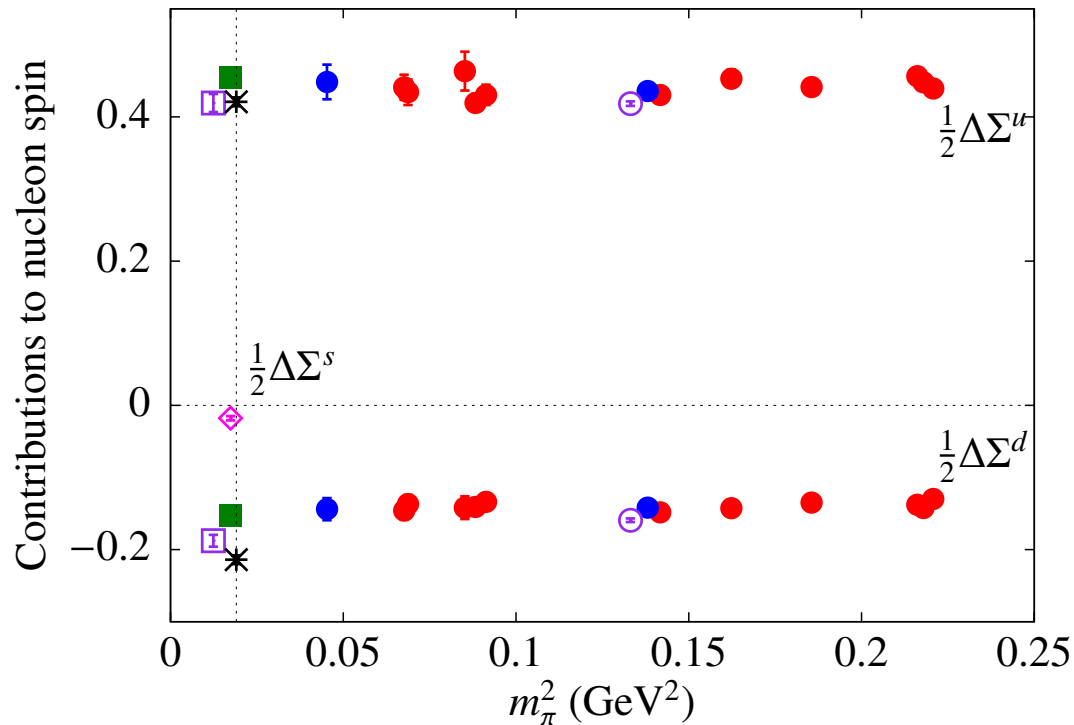
- Contribution exclusively by disconnected quark loop
- Results also at physical light quark masses
- Weak dependence on light quark mass

Contributions to nucleon spin

$$\Delta\Sigma^q = g_A^q$$

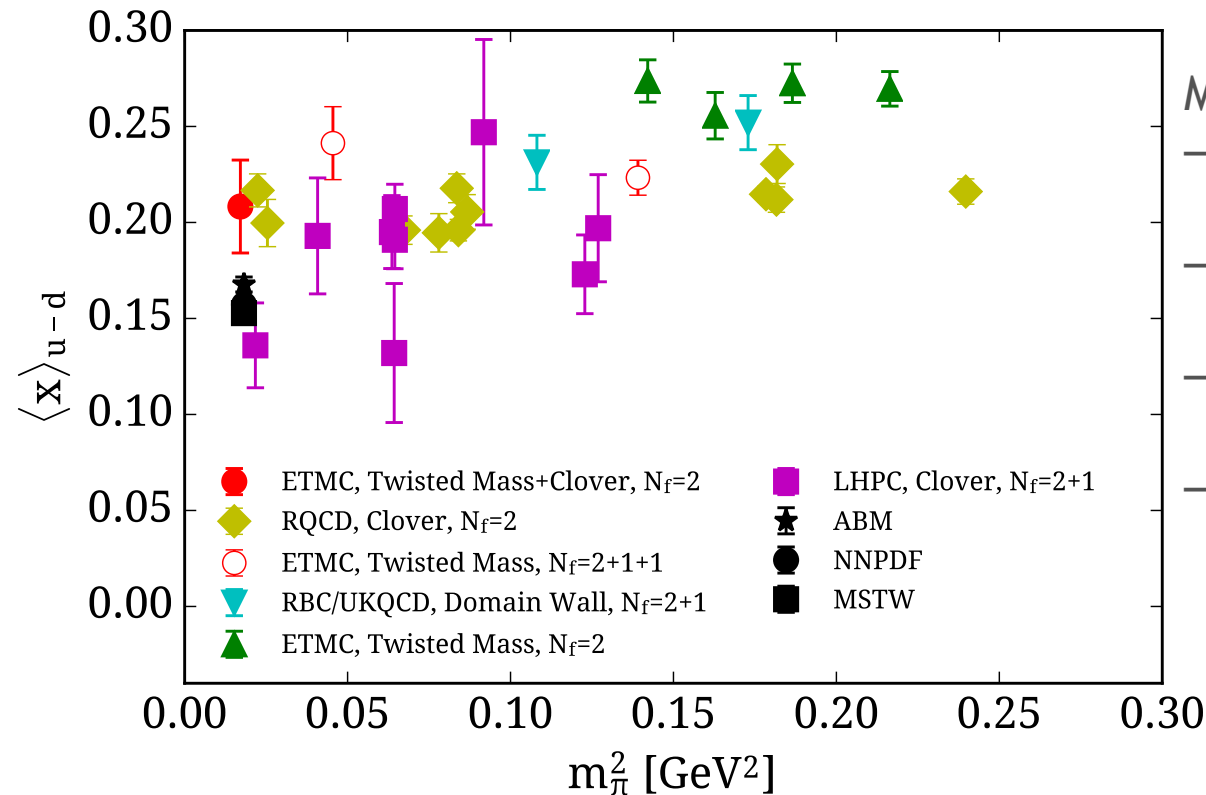
Axial charge – intrinsic quark spin contributions to nucleon spin

- Negative down-quark contribution
- Small strange-quark contribution
- Including disconnected contributions (open symbols) – better agreement with experiment (asterisks)



Momentum fraction

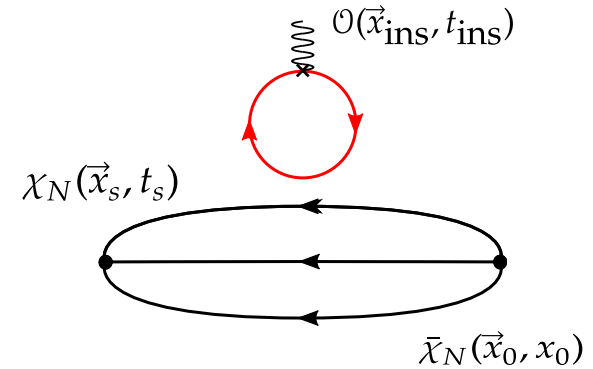
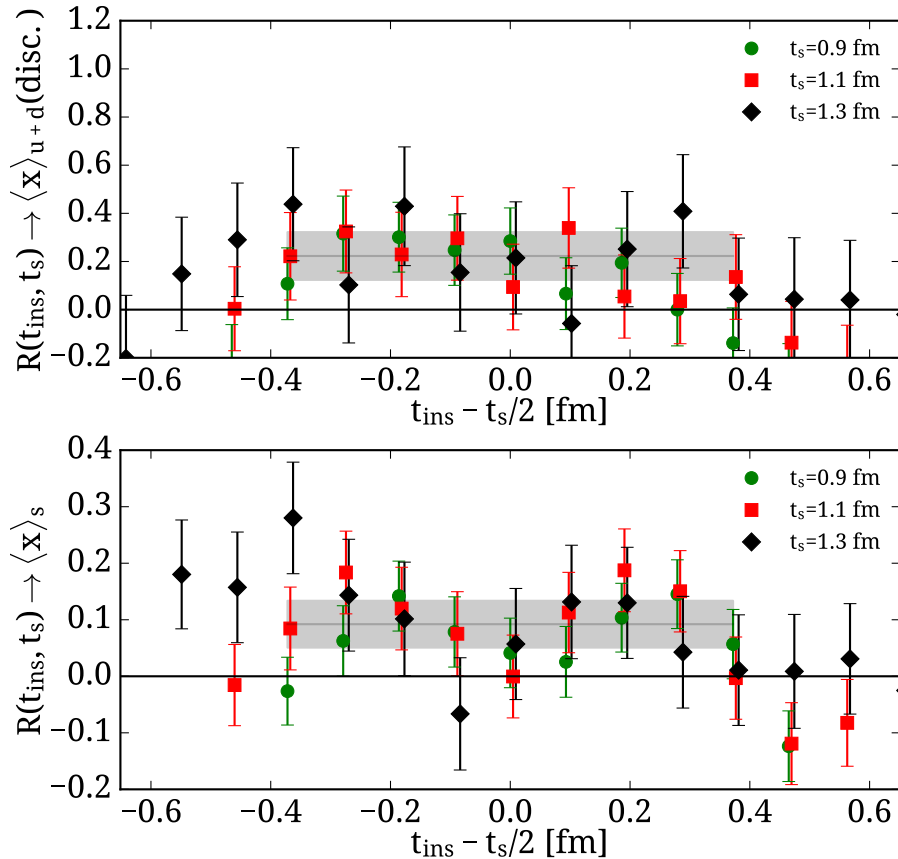
$$\mathcal{O}_V^{\mu\nu} = \bar{\psi} \gamma^{\{\mu} i D^{\nu\}} \psi \rightarrow \langle x \rangle_q$$



Momentum fraction

- General agreement between lattice actions
- More recent results at the physical point
- Pion mass dependence towards experimental values
- Need large separations ~1.5 fm

Momentum fraction



Disconnected contribution

- Needed for isoscalar
- $O(10^5)$ statistics to obtain statistically non-zero value
- Needed for individual quark-contributions
- To obtain spin-decomposition of nucleon

Spin decomposition

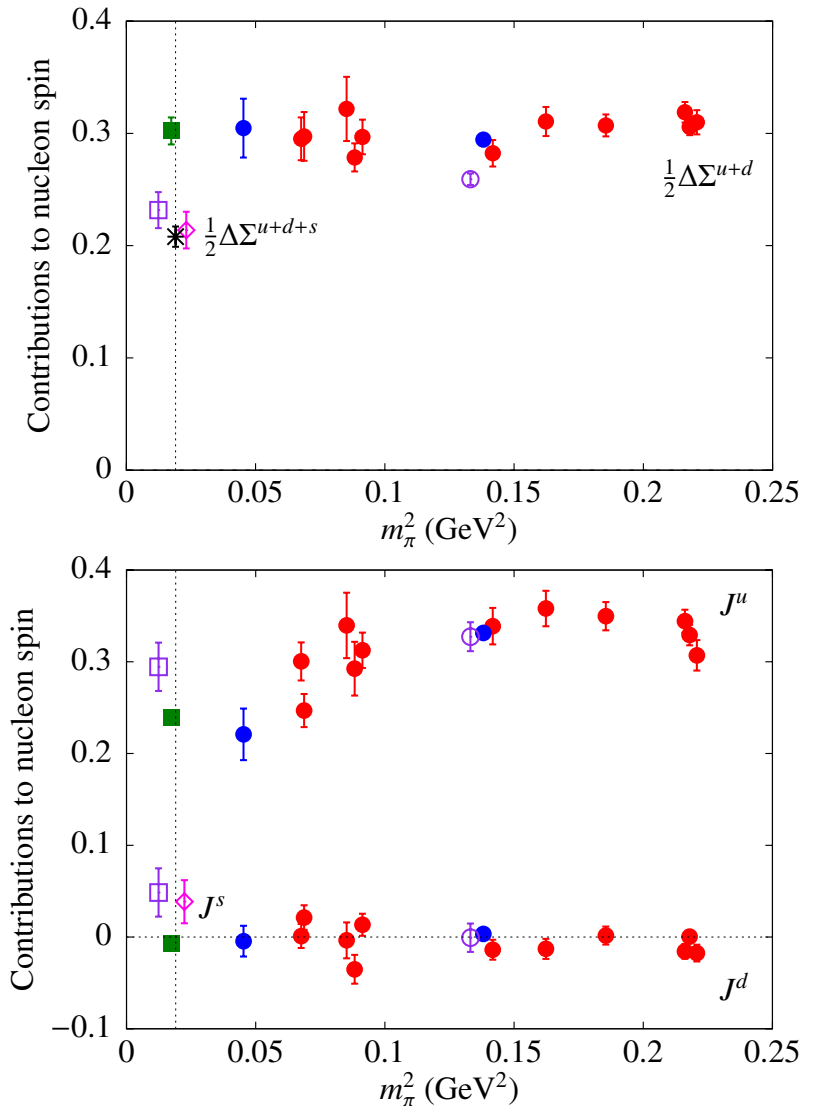
Ji's spin sum:
$$\frac{1}{2} = \sum_q \left(\frac{1}{2} \Delta\Sigma^q + L^q \right) + J^G$$

$$J^q = \frac{1}{2} \Delta\Sigma^q + L^q = \frac{1}{2} [A_{20}^q(0) + B_{20}^q(0)]$$

$$\Delta\Sigma^q = g_A^q$$

Nucleon spin

- Total J spin contributions from unpolarised matrix element A_{20} , B_{20} .
- Intrinsic spin contributions from axial matrix element $\Delta\Sigma$
- Orbital angular momentum of quarks from difference
- $J^{u+d+s} = 0.398(60)$, implies $J^G = 0.102(60)$



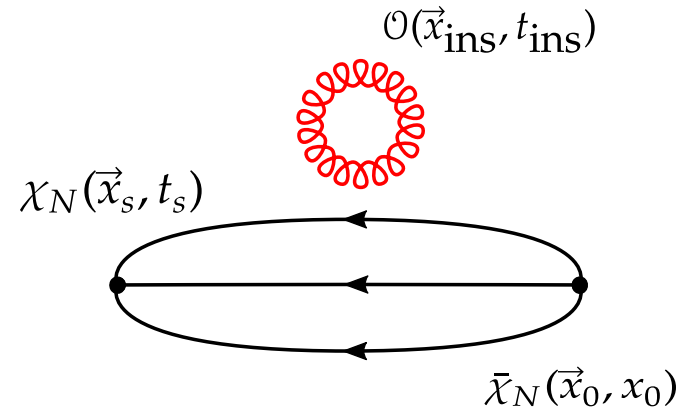
Gluon moment

Ji's spin sum:

$$\frac{1}{2} = \sum_q \left(\frac{1}{2} \Delta \Sigma^q + L^q \right) + J^G$$

$$J^G = \frac{1}{2} [A_{20}^G(0) + B_{20}^G(0)]$$

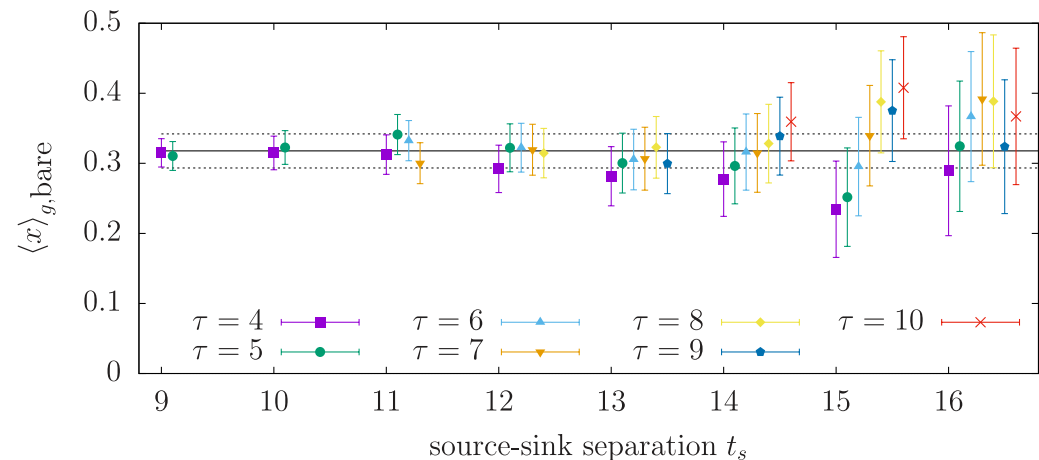
$$\mathcal{O} = \frac{2}{9} \frac{\beta}{a^4} \left[\sum_i \Re(P_{i4}) - \sum_{i < j} \Re(P_{ij}) \right]$$



Renormalisation

- Mixing with quark operator

$$\begin{pmatrix} J_q \\ J_G \end{pmatrix} = \begin{pmatrix} Z_{qq} & Z_{qG} \\ Z_{Gq} & Z_{GG} \end{pmatrix} \begin{pmatrix} J_q^{\text{bare}} \\ J_G^{\text{bare}} \end{pmatrix}$$



Spin decomposition

Ji's spin sum: $\frac{1}{2} = \sum_q \left(\frac{1}{2} \Delta \Sigma^q + L^q \right) + J^G$

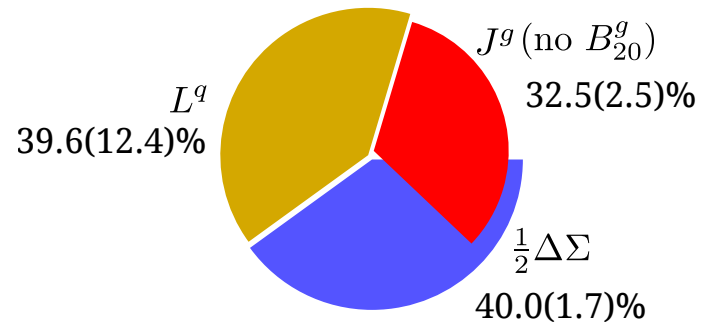
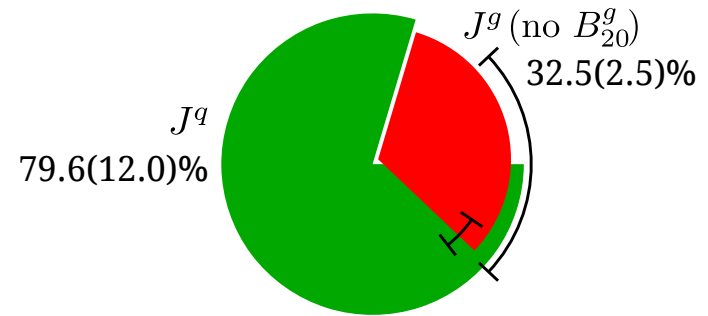
$$J^q = \frac{1}{2} \Delta \Sigma^q + L^q = \frac{1}{2} [A_{20}^q(0) + B_{20}^q(0)]$$

$$\Delta \Sigma^q = g_A^q$$

$$J^G = \frac{1}{2} [A_{20}^G(0) + B_{20}^G(0)]$$

Nucleon spin

- $J^G = 0.160(12)$ from direct calculation
PoS(DIS2016) [arXiv:1609.00253v1]
- $J^{u+d+s+G} = 0.541(59)$



Summary and outlook

★Lattice QCD in new era

- Physical pion mass simulations from a number of collaborations
- Other systematic uncertainties coming under control

★Nucleon spin

- Axial charge and momentum fraction at the physical point
- Results for disconnected fermion loops thanks to improved methods, order of magnitude more statistics
- New gluon moment calculations reveal gluon contribution to spin

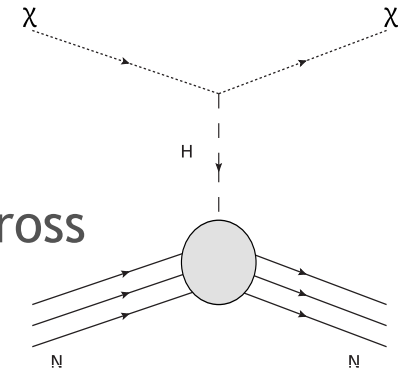
★What to expect

- Thorough estimation of systematic errors: continuum and infinite volume limits
- B_{20} calculation for disconnected and gluon
- Longer term: effects of isospin breaking and electromagnetic effects

Backup slides

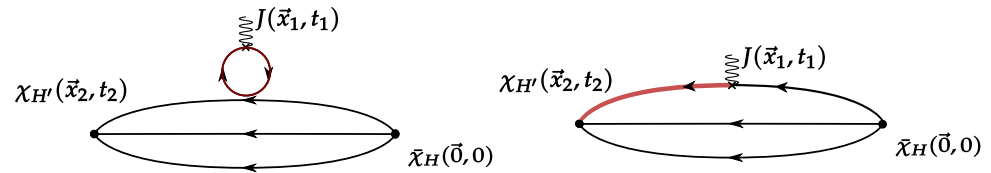
Nucleon sigma – terms

- Pion nucleon σ -term: $\sigma_{\pi N} = m_{ud} \langle N | \bar{u}u + \bar{d}d | N \rangle$
- Strange σ -term: $\sigma_s = m_s \langle N | \bar{s}s | N \rangle$
- Enter super-symmetric candidate particle scattering cross sections with nucleon (e.g. neutralino through Higgs)

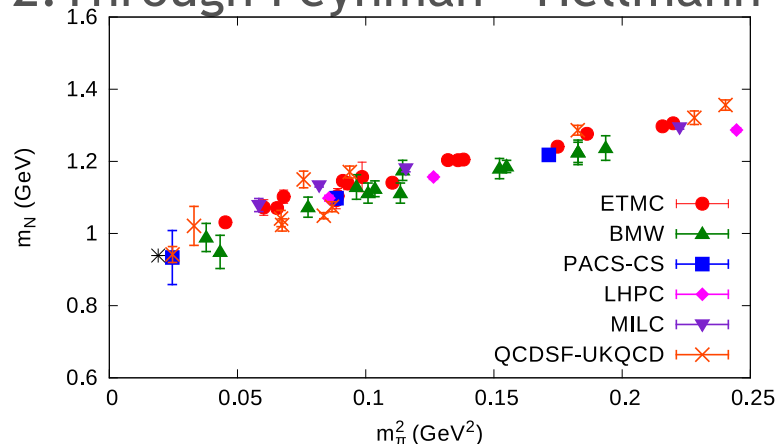


1. Direct calculation of matrix elements

Involves disconnected contributions

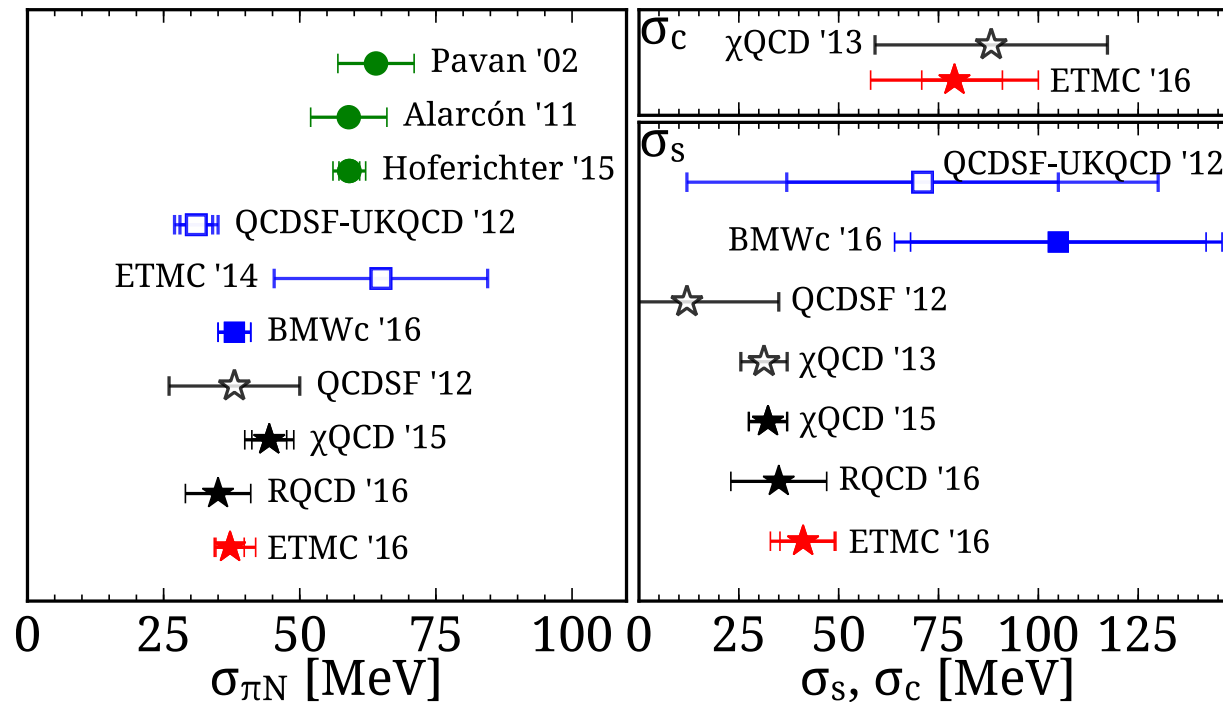


2. Through Feynman - Hellmann theorem: $\sigma_{\pi N} = m_{ud} \frac{\partial m_N}{\partial m_{ud}}$ $\sigma_s = m_s \frac{\partial m_N}{\partial m_s}$



- Reliance on effective theories for dependence on m_π
- Weak dependence on m_s

Nucleon sigma – terms

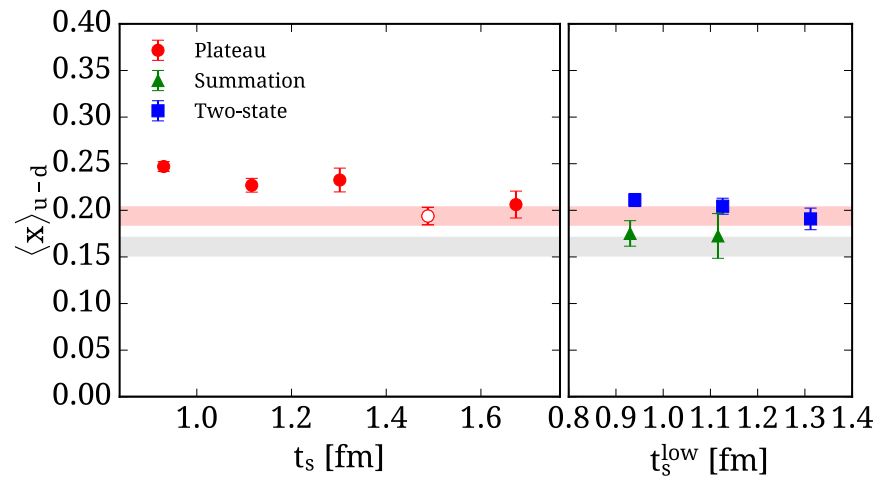
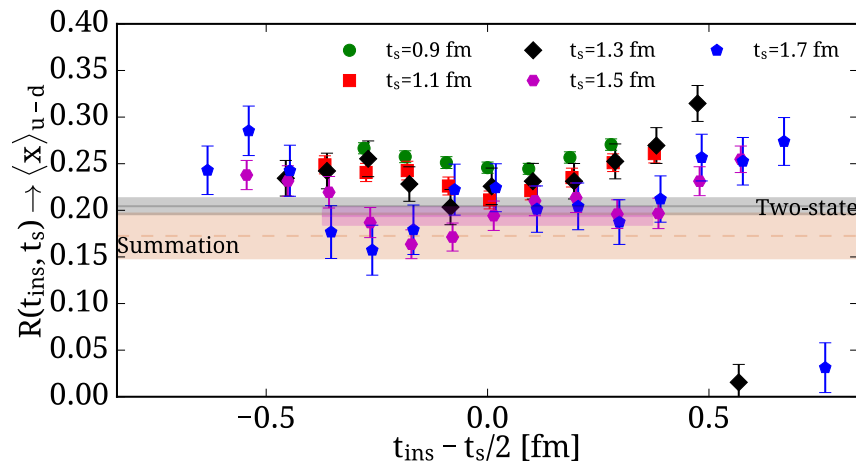


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- Recent results using direct matrix element evaluation [★]
- Large errors in FH-method [■](especially strange) due to sensitivity to quark mass
- Compare to phenomenology [●]

Momentum fraction

- Large sink-source separation
- >50000 statistics



Convergence of strange and light disconnected axial charge

